

CONVERGENCE ACCELERATION FOR FOURIER SERIES

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The paper studies convergence acceleration for Fourier series based on Pade approximants. An application to expansions by eigenfunctions of a boundary value problem for a first order model differential equation with discontinuous coefficient is considered and numerical results discussed.

§1. INTRODUCTION

It is well known that approximation of a 2-periodic $f \in C^\infty(R)$ function by truncated Fourier series (partial sum)

$$S_N(f) := \sum_{n=-N}^N f_n e^{i\pi n x}, \quad f_n = \frac{1}{2} \int_{-1}^1 f(x) e^{-i\pi n x} dx \quad (1)$$

is highly effective. When the approximated function has a discontinuity point, this truncation procedure leads to the Gibbs phenomena. To counter them, different solutions have been suggested in the literature (see [2], [12, 13] and the references therein). Thus A. Krylov in 1906 [14] and in 1966 by Lanczos [15] suggested subtracting a polynomial representing the discontinuities of the function and some of its derivatives. In [15] the correction polynomial was a linear combination of Bernoulli polynomials. In a series of papers [3, 5 – 8, 10, 11] Gottlieb and Eckhoff developed this method for practical realizations. Further we refer to this as Polynomial (or P-) method. Another way suggested in a general form by Cheney [4], is Fourier-Pade approximation which uses Pade approximants [1]. Other trigonometric-rational investigations were carried out in [9], [16].

In [17, 19] and [20] Pade approximants were applied to asymptotic expansion of Fourier coefficients. This approach leads to quasipolynomial approximation (QP-method) and actually generalizes the

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P-method.

The present paper applies Fourier-Pade approximants to P and QP approximations for additional acceleration of convergence and applies that approach to expansions by the eigenfunctions of a model problem for some differential equation (see [21]) with a non-smooth coefficients.

§2. POLYNOMIAL-PADE (PP-) APPROXIMATIONS

First we describe P-approximation. Suppose $f \in C^q[-1, 1]$, $q \geq 0$ and denote

$$A_k(f) = f^{(k)}(1) - f^{(k)}(-1), \quad k = 0, \dots, q.$$

According to asymptotic expansion of Fourier coefficients

$$f_n = \frac{(-1)^{n+1}}{2} \sum_{k=0}^{q-1} \frac{A_k(f)}{(i\pi n)^{k+1}} + \frac{1}{2(i\pi n)^q} \int_{-1}^1 f^{(q)}(x) e^{-i\pi n x} dx \quad (2)$$

the function f can be split into two parts

$$f(x) = F(x) + \sum_{k=0}^{q-1} A_k(f) B_k(x), \quad (3)$$

where F is a relatively smooth function with Fourier coefficients $F_n = o(n^{-q})$, $n \rightarrow \infty$ and $B_k(x)$ are 2-periodic Bernoulli polynomials with Fourier coefficients

$$B_{k,n} = \begin{cases} 0, & n = 0 \\ \frac{(-1)^{n+1}}{2(i\pi n)^{k+1}}, & n = \pm 1, \pm 2, \dots \end{cases}$$

Approximation of F by $S_N(F)$ leads to *Polynomial (P-) approximation*

$$S_{q,N}(f) = S_N(F) + \sum_{k=0}^{q-1} A_k(f) B_k(x), \quad (4)$$

where the Fourier coefficients of F can be found from (3)

$$F_n = f_n - \sum_{k=0}^{q-1} A_k(f) B_{k,n}.$$

We will write $S_{0,N}(f) \equiv S_N(f)$.

Now we apply Fourier-Pade approximation for additional acceleration of $S_N(F)$. Following [16] we consider a finite sequence of complex numbers $\theta := \{\theta_k\}_{|k|=1}^p$, $p \geq 1$ and denote

$$\begin{aligned} \Delta_n^0(\theta, F) &= F_n, \\ \Delta_n^k(\theta, F) &= \Delta_n^{k-1}(\theta, F) + \theta_k \operatorname{sgn}(n) \Delta_{(|n|-1)\operatorname{sgn}(n)}^{k-1}(\theta, F), \quad k \geq 1, \end{aligned}$$

where $\text{sgn}(n) = 1$ if $n \geq 0$ and $\text{sgn}(n) = -1$ if $n < 0$.

From (3), (4), we get

$$\begin{aligned} R_{q,N}(f) &:= f(x) - S_{q,N}(f) = F(x) - S_N(F) = R_N^+(F) + R_N^-(F), \\ R_N^+(F) &:= \sum_{n=N+1}^{\infty} F_n e^{i\pi n x}, \quad R_N^-(F) := \sum_{n=-\infty}^{-N-1} F_n e^{i\pi n x}. \end{aligned} \quad (5)$$

It is easy to check that

$$R_N^+(F) = -\frac{\theta_1 F_N e^{i\pi(N+1)x}}{1 + \theta_1 e^{i\pi x}} + \frac{1}{1 + \theta_1 e^{i\pi x}} \sum_{n=N+1}^{\infty} \Delta_n^1(\theta, F) e^{i\pi n x}.$$

Reiteration of this transformation p times leads to the expansion

$$\begin{aligned} R_N^+(F) &= -e^{i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_k \Delta_N^{k-1}(\theta, F)}{\prod_{s=1}^k (1 + \theta_s e^{i\pi x})} + \\ &+ \frac{1}{\prod_{k=1}^p (1 + \theta_k e^{i\pi x})} \sum_{n=N+1}^{\infty} \Delta_n^p(\theta, F) e^{i\pi n x}. \end{aligned} \quad (6)$$

Similar expansion of $R_N^-(F)$ reduces to the following *Polynomial-Pade (PP-) approximation* [16]

$$\begin{aligned} S_{p,q,N}(f) &:= S_N(F) - e^{i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_k \Delta_N^{k-1}(\theta, F)}{\prod_{s=1}^k (1 + \theta_s e^{i\pi x})} - \\ &- e^{-i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_{-k} \Delta_{-N}^{k-1}(\theta, F)}{\prod_{s=1}^k (1 + \theta_{-s} e^{-i\pi x})} + \sum_{k=0}^{q-1} A_k(f) B_k(x). \end{aligned}$$

It is natural to put $S_{0,q,N}(f) \equiv S_{q,N}(f)$.

There are different ways for determination of the vector θ . One option is Fourier-Pade method where the vector θ is found as a solution of the system

$$\Delta_n^p(\theta, F) = 0, \quad n = -N - p, \dots, -N - 1, N + 1, \dots, N + p.$$

Another option is connected with the following theorem, where $\|f\| = \left(\int_{-1}^1 |f(x)|^2 dx \right)^{1/2}$ denotes the L_2 -norm.

Theorem 1 [16]. Suppose $f \in C^{q+p}[-1, 1]$, for some $q \geq 0$, $p \geq 1$, and $f^{(q+p)}$ is absolutely continuous on $[-1, 1]$. If

$$\theta_k = \theta_{-k} = 1 - \frac{\tau_k}{N}, \quad k = 1, \dots, p, \quad \tau_k > 0, \quad \tau_j \neq \tau_i, \quad j \neq i;$$

then

$$\lim_{N \rightarrow \infty} N^{q+\frac{1}{2}} \|f(x) - S_{p,q,N}(f)\| = |A_q(f)|c_p(q),$$

where

$$c_p(q) = \frac{1}{\pi^{q+1}} \left(\int_1^\infty |\phi_{p,q}(t)|^2 dt \right)^{1/2},$$

$$\phi_{p,q}(t) := \frac{(-1)^p}{t^{q+1}} - \frac{1}{q!} \sum_{j=1}^p \frac{e^{-\tau_j(t-1)}}{\prod_{i \neq j}^p (\tau_i - \tau_j)} \sum_{k=0}^p \gamma_k(p) (-1)^{k+1} \sum_{m=0}^{p-k-1} (q+p-k-m-1)! \tau_j^m$$

and $\gamma_k(p)$ are defined by the identity

$$\prod_{k=1}^p (1 + \tau_k x) \equiv \sum_{k=0}^p \gamma_k(p) x^k. \quad (7)$$

It can be easily shown that for $p = 0$ and $f \in C^q[-1, 1]$, for some $q \geq 0$, with absolutely continuous q -th derivative

$$\lim_{N \rightarrow \infty} N^{q+\frac{1}{2}} \|f(x) - S_{q,N}(f)\| = |A_q(f)|c_0(q), \quad c_0(q) = \frac{1}{\pi^{q+1} \sqrt{2q+1}}.$$

In Table 1 we represent some results from [16] on the choice of parameters τ_k that minimize the L_2 -error for $p = 3$. The ratio $c_0(q)/c_3(q)$ describes effectiveness of L_2 -optimal rational approximation $S_{p,q,N}(f)$ compared to $S_{q,N}(f)$ for $N \gg 1$.

q	1	2	3	4	5	6
$c_3(q)$	0.00095	0.00007	$9 \cdot 10^{-6}$	$1 \cdot 10^{-6}$	$2 \cdot 10^{-7}$	$4 \cdot 10^{-8}$
$c_0(q)/c_3(q)$	61.3	185.1	411.6	771.8	1296.7	2017.4
τ_1	0.2510	0.6382	1.1230	1.6730	2.2699	2.9023
τ_2	1.28553	2.2362	3.2067	4.1868	5.1725	6.1617
τ_3	4.2225	5.7813	7.2573	8.6781	10.0589	11.4089

Table 1. Numerical values of $c_3(q)$ and $c_0(q)/c_3(q)$ for $1 \leq q \leq 6$ using the numerical optimal values of parameters τ_k , $k = 1, 2, 3$.

§3. QUASIPOLYNOMIAL-PADE (QPP-) APPROXIMATIONS

Following [17 – 20] we consider a finite sequence of complex numbers $\eta := \{\eta_k\}_{k=1}^m$, $m \geq 1$ and denote

$$\delta_n^0(\eta, f) = A_n(f), \quad \delta_n^k(\eta, f) = \delta_n^{k-1}(\eta, f) + \eta_k \delta_{n-1}^{k-1}(\eta, f), \quad 1 \leq k \leq q.$$

If $n < 0$, we put $\delta_n^k(\eta, f) = 0$, $k = 0, 1, \dots$

It can easily be checked that

$$\sum_{k=0}^{q-1} A_k(f)x^k = x^q \frac{A_{q-1}(f)\eta_1}{1 + \eta_1 x} + \frac{1}{1 + \eta_1 x} \sum_{k=0}^{q-1} (A_k(f) + \eta_1 A_{k-1}(f))x^k. \quad (8)$$

Note that for $\eta_1 = 0$ the sum in the left side of (8) remains unchanged. Reiteration of this transformation m times ($m \leq q-1$) leads to the formula

$$\sum_{k=0}^{q-1} A_k(f)x^k = x^q \sum_{k=1}^m \frac{\eta_k \delta_{q-1}^{k-1}(\eta, f)}{\prod_{s=1}^k (1 + \eta_s x)} + \frac{1}{\prod_{s=1}^m (1 + \eta_s x)} \sum_{k=0}^{q-1} \delta_k^m(\eta, f)x^k. \quad (9)$$

Now suppose $f \in C^q[-1, 1]$ for some $q \geq 1$. Applying transformation (9) to the first term of (2) with $(i\pi n)^{-1}$ instead of x we get

$$f_n = P_n + Q_n, \quad n \neq 0, \quad (10)$$

where

$$\begin{aligned} P_n &= \frac{(-1)^{n+1}}{2(i\pi n)^{q+1}} \sum_{k=1}^m \frac{\eta_k \delta_{q-1}^{k-1}(\eta, f)(i\pi n)^k}{\prod_{s=1}^k (i\pi n + \eta_s)} \\ &+ \frac{(-1)^{n+1}(i\pi n)^m}{2 \prod_{k=1}^m (i\pi n + \eta_k)} \sum_{k=q-m}^{q-1} \frac{\delta_k^m(\eta, f)}{(i\pi n)^{k+1}} + \frac{1}{2(i\pi n)^q} \int_{-1}^1 f^{(q)}(t) e^{-i\pi n t} dt. \end{aligned} \quad (11)$$

and

$$Q_n = \frac{(-1)^{n+1}(i\pi n)^m}{2 \prod_{s=1}^m (i\pi n + \eta_s)} \sum_{k=0}^{q-m-1} \frac{\delta_k^m(\eta, f)}{(i\pi n)^{k+1}}.$$

According to (10) the function f can be split into two parts

$$f(x) = P(x) + Q(x), \quad (12)$$

where

$$P(x) = \sum_{n=-\infty}^{\infty} P_n e^{i\pi n x}, \quad P_0 = f_0, \quad Q(x) = \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} Q_n e^{i\pi n x}.$$

Approximation of P by the truncated Fourier series leads to the following approximation [17 – 20]

$$T_{q,m,N}(f) = S_N(P) + Q(x), \quad (13)$$

where the Fourier coefficients of P can be found from (12)

$$P_n = f_n - Q_n.$$

The unknown vector η in (13) we determine from the system

$$\delta_k^m(\eta, f) = 0, \quad k = q - m, \dots, q - 1. \quad (14)$$

In [17] it was shown that the function $Q(x)$ is a quasipolynomial of the form

$$Q(x) = \sum_k a_k x^{p_k} e^{i\omega_k x},$$

where $\omega_k \in \mathbb{C}$ and $\{p_k\}$ is a set of nonnegative integers. The approximation (13), (14) we call *QP-method or QP-approximation*. It is important to note that for $\eta_1 = \eta_2 = \dots = \eta_m = 0$ the QP-approximation coincides with P-approximation $S_{q,N}(f)$.

For additional acceleration of QP-method we proceed as in the previous section. From (13) and (12), we have, see (5)

$$R_{q,m,N}(f) := f(x) - T_{q,m,N}(f) = R_N^+(P) + R_N^-(P),$$

where

$$\begin{aligned} R_N^+(P) &:= \sum_{n=N+1}^{\infty} P_n e^{i\pi n x} \\ &= -e^{i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_k \Delta_N^{k-1}(\theta, P)}{\prod_{s=1}^k (1 + \theta_s e^{i\pi s x})} + \frac{1}{\prod_{k=1}^p (1 + \theta_k e^{i\pi k x})} \sum_{n=N+1}^{\infty} \Delta_n^p(\theta, P) e^{i\pi n x}. \end{aligned}$$

Similar expansion of $R_N^-(P)$ reduces to the following *Quasipolynomial-Pade (QPP-) approximation*:

$$\begin{aligned} T_{p,q,m,N}(f) &= S_N(P) - e^{i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_k \Delta_N^{k-1}(\theta, P)}{\prod_{s=1}^k (1 + \theta_s e^{i\pi s x})} \\ &\quad - e^{-i\pi(N+1)x} \sum_{k=1}^p \frac{\theta_{-k} \Delta_{-N}^{k-1}(\theta, P)}{\prod_{s=1}^k (1 + \theta_{-s} e^{-i\pi s x})} + Q(x). \end{aligned}$$

Now we prove an analog of Theorem 1 for QPP-approximations.

In [16] the following lemma was proved:

Lemma 1 [16]. Suppose that the sequence P_n has the following asymptotic expansion for $q \geq 0$, $p \geq 1$:

$$P_n = \frac{(-1)^{n+1}}{2} \sum_{s=q}^{p+q} \frac{\alpha_s}{(i\pi n)^{s+1}} + o(n^{-p-q-1}), \quad n \rightarrow \infty,$$

where $\{\alpha_s\}_{s=q}^{p+q}$ are some constants. If

$$\theta_k = \theta_{-k} = 1 - \frac{\tau_k}{N}, \quad k = 1, \dots, p,$$

then the asymptotic expansion

$$\Delta_n^p(\theta, P) = \alpha_q \frac{(-1)^{n+p+1}}{2(i\pi)^{q+1}q!} \sum_{k=0}^p \frac{(q+p-k)!(-1)^k \gamma_k(p)}{N^k(n-k)^{q+1}|n-k|^{p-k}} + o(n^{-q-p-1}),$$

holds as $N \rightarrow \infty$, $|n| \geq N+1$, where the numbers $\gamma_k(p)$ are defined by (7).

By $\mu_k(m), k = 0, \dots, m$, we denote the coefficients of the polynomial

$$\prod_{k=1}^m (1 + \eta_k x) \equiv \sum_{k=0}^m \mu_k(m) x^k.$$

Note that the system (14) can be written in the following form:

$$\sum_{s=1}^m \mu_s(m) A_{k-s+q-m-1}(f) = -A_{k+q-m-1}(f), \quad k = 1, \dots, m, \quad (15)$$

and denote

$$U_r^m = [A_{k-s+r}(f)], \quad k, \quad s = 1, \dots, m.$$

Theorem 2. Suppose $f \in C^{q+p}[-1, 1]$ for some $q \geq 1$, $p \geq 0$ and $f^{(q+p)}$ is absolutely continuous on $[-1, 1]$.

If $\det U_{q-m-1}^m \neq 0$, then with η found from (14) and $\theta_k = \theta_{-k} = 1 - \frac{\tau_k}{N}$, $k = 1, \dots, p$, $\tau_k > 0$, $\tau_j \neq \tau_i$, $j \neq i$,

$$\lim_{N \rightarrow \infty} N^{q+\frac{1}{2}} \|f - T_{p,q,m,N}(f)\| = \left| \frac{\det U_{q-m}^{m+1}}{\det U_{q-m-1}^m} \right| c_p(q),$$

where $c_p(q)$ is defined in theorem 1.

Proof: From (12), (13) and (6) we have

$$f(x) - T_{p,q,m,N}(f) := R_{p,q,m,N}^+(P) + R_{p,q,m,N}^-(P),$$

where

$$R_{p,q,m,N}^\pm(P) = \frac{1}{\prod_{k=1}^p (1 + \theta_{\pm k} e^{\pm i\pi x})} \sum_{n=N+1}^{\infty} \Delta_{\pm n}^p(\theta, P) e^{\pm i\pi n x}.$$

After simple calculations we obtain from (11)

$$P_n = \frac{(-1)^{n+1}}{2} \sum_{\ell=q}^{p+q} \frac{\alpha_\ell}{(i\pi n)^{\ell+1}} + o(n^{-p-q-1}), \quad n \rightarrow \infty,$$

where

$$\alpha_\ell = A_\ell(f) + \sum_{k=1}^m \eta_k \delta_{q-1}^{k-1}(\eta, f) \sum_{s=1}^k \frac{\eta_s^{k-1} (-1)^{\ell-q} \eta_s^{\ell-q}}{\prod_{\substack{j=1 \\ j \neq s}}^k (\eta_s - \eta_j)}.$$

Now we can apply Lemma 1 to the sequence P_n taking into account that

$$\alpha_q = A_q(f) + \sum_{k=1}^m \eta_k \delta_{q-1}^{k-1}(\eta, f) = \delta_q^m(\eta, f),$$

which follows from the relation

$$\begin{aligned} \delta_q^m(\eta, f) &= \delta_q^{m-1}(\eta, f) + \eta_m \delta_{q-1}^{m-1}(\eta, f) = \delta_q^{m-2}(\eta, f) + \eta_{m-1} \delta_{q-1}^{m-2}(\eta, f) + \eta_m \delta_{q-1}^{m-1}(\eta, f) \\ &= \delta_q^0(\eta, f) + \sum_{k=1}^m \eta_k \delta_{q-1}^{k-1}(\eta, f) = A_q(f) + \sum_{k=1}^m \eta_k \delta_{q-1}^{k-1}(\eta, f). \end{aligned}$$

By Cramer's rule, from (15) we get

$$\mu_s(m) = \frac{M_s}{\det U_{q-m-1}^m}, \quad s = 1, \dots, m,$$

where $\{M_s\}$ are the corresponding minors. Consequently,

$$\begin{aligned} \delta_q^m(\eta, f) &= A_q(f) + \sum_{s=1}^m \mu_s(m) A_{q-s}(f) = A_q(f) + \frac{1}{\det U_{q-m-1}^m} \sum_{s=1}^m M_s A_{q-s}(f) \\ &= (-1)^m \frac{\det U_{q-m}^{m+1}}{\det U_{q-m-1}^m}. \end{aligned}$$

To get the proof it remains to proceed as in the proof of Theorem 1 with $A_q(f)$ replaced by

$$(-1)^m \frac{\det U_{q-m}^{m+1}}{\det U_{q-m-1}^m}.$$

Note that for $p = 0$ Theorem 2 was proved in [20].

§4. NUMERICAL RESULTS

For any given f , q and m we put

$$a_{q,m}(f) = \left| A_q(f) \frac{\det (U_{q-m-1}^m)}{\det (U_{q-m}^{m+1})} \right|.$$

The constant $a_{q,m}(f)$ describes the effectiveness of QPP-approximation compared to PP-approximation (with the same value of parameter p), assuming $N \gg 1$ and Theorems 1,2 valid. Calculations show that this constant in fact describes the convergence acceleration in broader situations. Note that $a_{q,m}$ is independent of the parameter p .

Let us investigate the example

$$f(x) = \frac{\sin(5x - 0.2)}{1.1 - x}. \quad (16)$$

Figure 1 represents the graphs of $a_{q,m}(f)$ for (16) for $q = 5, 6, 7$ and $1 \leq m \leq q - 1$. For this function the QPP-method proves to be more precise than PP-method (for the same value of parameter p) almost 15 times for $q = 5; m = 3$ and 50 times when $q = 7; m = 3$.

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Figure 1: Graphics of $a_{q,m}(f)$ for (16) for $q = 5, 6, 7$ and $1 \leq m \leq q - 1$.

The relative effectiveness of QPP-method against PP-method can be described by the fraction

$$a_{N,q,m,p}(f) = \frac{\max_{|x| \leq 1} |f - S_{p,q,N}(f)|}{\max_{|x| \leq 1} |f - T_{p,q,m,N}(f)|}.$$

Table 1 shows approximate values of $a_{N,7,3,3}$ for (16). Calculations are carried out with 64 digits of precision by MATHEMATICA package.

N	32	64	128	256	512
$a_{N,7,3,3}$	101.17	64.36	50.91	48.58	48.57

Table 1: Approximate values of $a_{N,7,3,3}$ for different N .

Comparison with theoretical value $a_{7,3} = 49.4408$ shows that experimental and theoretical estimates are rather close for $N \geq 64$.

	$S_{0,q,N}$	$S_{1,q,N}$	$S_{2,q,N}$	$S_{3,q,N}$
128	1.8×10^{-9}	2.1×10^{-10}	4.5×10^{-11}	1.3×10^{-11}
256	7.3×10^{-12}	8.3×10^{-13}	1.7×10^{-13}	4.7×10^{-14}
	$T_{0,q,m,N}$	$T_{1,q,m,N}$	$T_{2,q,m,N}$	$T_{3,q,m,N}$
128	3.5×10^{-11}	4.1×10^{-12}	8.8×10^{-13}	2.5×10^{-13}
256	1.52×10^{-13}	1.7×10^{-14}	3.5×10^{-15}	9.7×10^{-16}

Table 2: Uniform errors by QPP and PP methods for different values of N, p and $q = 7, m = 3$.

Figure 1 shows also the optimal values of m when parameter q is fixed. Thus we see that for $q = 7$ the optimal is $m = 3$. Table 2 presents uniform errors in approximation of (16) by QPP and PP methods for different values of N, p and $q = 7, m = 3$. Comparison shows that $S_{2,7,N}$ has the same precision as $T_{0,7,3,N}$ while $T_{1,7,3,N}$ is 3 times more precise than $S_{3,7,N}$.

All calculations are carried out by the package MATHEMATICA with 64 digits of precision.

§ 5. A MODEL PROBLEM

The approach of previous sections was generalized in [18] for expansions by eigenfunctions for one-dimensional boundary problems in the case where the coefficients of the equations are smooth. Here we consider a simple first order model differential equation with non-smooth coefficient. Some preliminary results were obtained in [21].

We pose the eigenvalue problem

$$i \frac{du}{dx} = \lambda \varepsilon(x) u(x), \quad x \in (-1, 1), \quad (17)$$

$$u(-1) = u(1), \quad (18)$$

where $\varepsilon(x) > \delta > 0$, $\varepsilon^{(q+1)}$ is a piecewise-continuous function in $[-1, 1]$ with potential points of discontinuity among $\alpha = \{\alpha_k\}_{k=0}^\mu$, $-1 = \alpha_0 < \alpha_1 < \alpha_2 < \dots < \alpha_{\mu-1} < \alpha_\mu = 1$. We denote this class of functions by $C_\alpha^{q+1}[-1, 1]$.

It is easy to calculate the eigenvalues $\{\lambda_n\}$ and eigenfunctions $\{\phi_n\}$ of the problem (17), (18)

$$\phi_n(x) = e^{-i\eta n \int_{-1}^x \varepsilon(t) dt}, \quad \lambda_n = -\pi \eta n, \quad n \in \mathbb{Z},$$

where

$$\eta = \frac{2\pi}{\int_{-1}^1 \varepsilon(t) dt}.$$

The system $\{\phi_n(x)\}_{n=-\infty}^\infty$ is orthogonal in the weighted space $L_2[(-1, 1), \varepsilon]$.

We consider now the formal series

$$f(x) = \sum_{n=-\infty}^\infty f_n \phi_n(x), \quad f_n = \frac{\eta}{2\pi} \int_{-1}^1 \varepsilon(x) f(x) \overline{\phi_n(x)} dx \quad (19)$$

and the approximation formula:

$$W_N(f) = \sum_{n=-N}^N f_n \phi_n(x). \quad (20)$$

Lemma 2. *If $f \in C_\alpha^{q+1}[-1, 1]$ then the following asymptotic expansion holds*

$$f_n = H_n + G_n, \quad (21)$$

$$H_n = \frac{\eta}{2\pi} \sum_{\ell=0}^{\mu-1} \sum_{k=0}^q \frac{(-1)^k A_k(f, \alpha_\ell)}{(i\eta n)^{k+1}} e^{i\eta n \int_{-1}^{\alpha_\ell} \varepsilon(t) dt}, \quad G_n = \frac{\eta}{2\pi} \frac{(-1)^{q+1}}{(i\eta n)^{q+1}} \int_{-1}^1 \varepsilon(x) g_{q+1}(x) \overline{\phi_n(x)} dx,$$

where

$$g_0(x) = f(x), \quad g_k(x) = \frac{g'_{k-1}(x)}{\varepsilon(x)},$$

$$A_k(f, -1) = g_k(1) - g_k(-1), \quad A_k(f, x) = g_k(x-0) - g_k(x+0).$$

Proof: We divide the integral in (19) into parts by the jump points

$$f_n = \frac{\eta}{2\pi} \int_{-1}^1 \varepsilon(x) f(x) \overline{\phi_n(x)} dt = \frac{\eta}{2\pi} \sum_{\ell=0}^{\mu-1} \int_{\alpha_\ell}^{\alpha_{\ell+1}} \varepsilon(x) f(x) \overline{\phi_n(x)} dx.$$

Integration by parts yields

$$\begin{aligned} \frac{\eta}{2\pi} \int_{\alpha_\ell}^{\alpha_{\ell+1}} \varepsilon(x) f(x) \overline{\phi_n(x)} dx &= \frac{\eta}{2\pi} \int_{\alpha_\ell}^{\alpha_{\ell+1}} f(x) \varepsilon(x) e^{i\eta n \int_{-1}^x \varepsilon(t) dt} dx = \\ &= \frac{g_0(x) e^{i\eta n \int_{-1}^x \varepsilon(t) dt}}{2i\pi n} \Big|_{\alpha_\ell}^{\alpha_{\ell+1}} - \frac{1}{2i\pi n} \int_{\alpha_\ell}^{\alpha_{\ell+1}} g_1(x) \varepsilon(x) e^{i\eta n \int_{-1}^x \varepsilon(t) dt} dx. \end{aligned} \quad (22)$$

From (22) we get

$$\frac{\eta}{2\pi} \int_{-1}^1 \varepsilon(x) f(x) \overline{\phi_n(x)} dt = \sum_{l=0}^{\mu-1} \frac{A_0(f, \alpha_l)}{2i\pi n} e^{i\eta n \int_{-1}^{\alpha_l} \varepsilon(t) dt} - \frac{1}{2i\pi n} \int_{-1}^1 \varepsilon(x) g_1(x) \overline{\phi_n(x)} dx.$$

Integration by parts $q-1$ times leads to (21). The proof is complete.

For convergence acceleration of (20) we use the idea of P-approximation. According to Lemma 2 we split the function f into two parts

$$f(x) = H(x) + G(x)$$

where $G(x)$ is smooth as compared with $H(x)$, and consider the following analog of P-method:

$$W_{q,N}(f) = H(x) + \sum_{n=-N}^N G_n \phi_n(x).$$

Theorem 3. *If $f \in C_\alpha^{q+1}[-1, 1]$, then*

$$f - W_{q,N}(f) = o(N^{-q}), \quad N \rightarrow \infty.$$

Proof: immediately follows from the formulas

$$\begin{aligned} f - W_{q,N}(f) &= G - W_N(G) = \sum_{|n|>N} G_n \phi_n(x), \\ G_n &= \frac{(-1)^q \eta}{2\pi (i\eta n)^{q+1}} \int_{-1}^1 g_{q+1}(x) \varepsilon(x) e^{i\eta n \int_{-1}^x \varepsilon(t) dt} dx, \end{aligned} \quad (23)$$

and from the fact that integral in the right hand side of (23) is $o(1)$, $n \rightarrow \infty$, according to the Riemann–Lebesgue Theorem.

The analogs of PP and QPP approximations can be constructed similarly. By $W_{q,p,N}(f)$ and $W_{q,p,m,N}(f)$ we denote the analogs of PP and correspondingly QPP approximations for (20).

Consider the function

$$f(x) = \frac{1}{1.1 - x}. \quad (24)$$

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Figure 2: Uniform errors in logarithm scale while approximating (24) by $W_{q,p,N}(f)$ and $W_{q,p,m,N}(f)$ for $q = 7$, $m = 3$, $p = 0, 2$ and ε as in (25).

In Figure 2, the uniform errors in logarithm scale are shown while approximating (24) by $W_{q,p,N}(f)$ and $W_{q,p,m,N}(f)$ for $q = 7$, $m = 3$, $p = 0, 2$ and

$$\varepsilon(x) = \begin{cases} 1, & x < 1/3 \\ 4, & x \geq 1/3. \end{cases} \quad (25)$$

We see that $W_{7,2,N}(f)$ and $W_{7,2,3,N}(f)$ are 4 to 10 times more precise (the difference is higher for greater values of the parameter N) compared to approximations $W_{7,0,N}(f)$ and $W_{7,0,3,N}(f)$ correspondingly. Approximations $W_{7,0,3,N}(f)$ and $W_{7,2,N}(f)$ show the same precision (see also remarks to Table 2).

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